

# Combined Effects of Hall and Ion-Slip Currents on Steady Free-Convective Flow of an Incompressible Viscous and Electrically Conducting Fluid with Heat and Mass Transfer over a Porous Flat Plate Embedded in Porous Medium

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**Abstract:** The present paper deals with the heat and mass transfer of a steady free-convective laminar flow of an incompressible viscous and electrically conducting fluid over a porous flat plate in the presence of transverse magnetic field, taking into account Hall and ion slip currents. The non-linear governing equations together with the boundary conditions are reduced to a system of non-linear ordinary differential equations by using similarity transformations. The system of non-linear ordinary differential equations are solved by shooting procedure using fourth order Runge-Kutta Method. The important flow characteristics, velocity, temperature and concentration profiles as well as the skin-friction coefficient  $C_f$ , local Nusselt number  $Nu$  and the local Sherwood number  $Sh$  for various values of the non-dimensional parameters like Magnetic parameter ( $M$ ), Hall current ( $\beta_e$ ), Ion-slip parameter ( $\beta_i$ ), Schmidt Number ( $Sc$ ), Soret Number ( $Sr$ ), Suction/Blowing parameter ( $S$ ) are studied numerically.

**Keywords -** Hall current, MHD, Ion-Slip current, Porous medium, Heat and mass transfer

## 1. Introduction

A few decades ago, many researchers paid more attention to MHD fluid flow problems but ignored the Hall current and Ion-slip current in Ohms law.

The MHD heat and mass transfer flow under the strong magnetic field plays a key role in many engineering problems such as astrophysics, chemical engineering and geophysics. In flows of laboratory plasma, Hall and ion-slip currents are likely to be important. Hence it is now proposed to study heat and mass transfer of a free convection MHD flow past an infinite vertical porous plate taking into account Hall and Ion slip currents. Combined effects of hall and ion-slip currents on free convective flow past a semi-infinite vertical plate was investigated by Abo-Eldahab *et al.* [1]. The effect of Hall and Ion-slip currents on fully developed electrically conducting fluid flow between vertical parallel plates in the presence of a temperature dependent heat source was investigated using homotopy analysis method by

Srinivasacharya and Kaladhar [2]. Later, Darbhasayanam Srinivasacharya and Kolla Kaladhar [3] considered the effect of Hall and Ion-slip currents on electrically conducting couple stress fluid flow between two circular cylinders in the presence of a temperature dependent heat source. Steady motion of electrically conducting viscous fluid in the presence of Hall current was studied by Tani [4]. The effects of Hall, ion-slip currents, and variable thermal diffusivity on magnetomicro-polar fluid flow, heat and mass transfer with suction through a porous medium was numerically analyzed by Motsa and Shateyi [5]. Attia [6] studied the flow of a dusty fluid in the presence of Hall and ion slip current using analytical procedure. Shampa Ghosh [10] studied magneto hydrodynamic boundary layer flow over a stretching sheet with chemical reaction.

The heat transfer of MHD Couette flow in the presence of Hall and Ion-slip currents was studied by Soundelgekar [7]. Unsteady MHD Couette flow of an incompressible viscous fluid in the presence of Hall and Ion-slip currents in a porous medium

was studied by Nirmal Ghara *et al.* [8]. The unsteady Couette flow of an electrically conducting viscous incompressible fluid between two parallel horizontal non-conducting porous plates with heat transfer in the presence of Ion-slip was studied by Attia [9].

The absence of Hall current and ion-slip current in Shampa Ghosh [10] motivated us to take up the present work wherein we study the effect of Hall current and ion-slip current on heat and mass transfer in a steady free-convective flow in the presence of a uniform transverse magnetic field in a porous medium. The governing equations are transformed by a similarity transformation into a system of non-linear ordinary differential equations which are solved numerically by fourth order Runge-Kutta Method. Numerical calculations are performed for various values of the magnetic parameter, Hall current parameter, Ion-slip parameter, Schimdt number and Soret number. The results are given for the coefficients of skin-friction, rate of heat transfer and rate of mass transfer for various values of magnetic, Hall current and Ion-slip parameters.

## 2. Mathematical Model

We consider the steady free-convective laminar flow of an incompressible viscous and electrically conducting fluid with heat and mass transfer over a porous flat plate in the presence of transverse magnetic field. Magnetic field is assumed to be strong enough to produce Hall and ion-slip currents. We choose a stationary frame of reference  $(x, y, z)$  such that x-axis is chosen along the direction of motion, y-axis is normal to the surface and z-axis is transverse to the xy-plane. For an electrically conducting fluid, the Hall and ion-slip currents affect the flow significantly in the presence of large magnetic field. The temperature and the concentration are maintained at prescribed constant values  $T_w, C_w$  at the plate and  $T_\infty$  and  $C_\infty$  are the fixed values far away from the sheet.

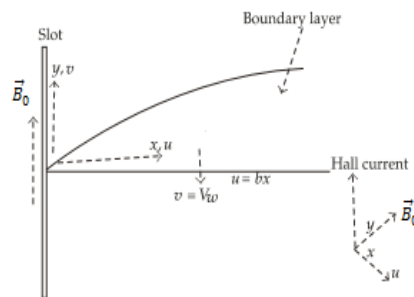


Fig 1 Schematic diagram of the problem

Owing to the above assumptions, the free-convection flow with heat and mass transfer and generalized Ohm's law with Hall current effect are governed by the following system of equations:

$$u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} - \frac{\nu}{k} u - \frac{B_0}{\rho} j_z \tag{2}$$

$$u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} = \nu \frac{\partial^2 w}{\partial y^2} - \frac{\nu}{k} w + \frac{B_0}{\rho} j_x \tag{3}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{\kappa}{\rho C_p} \frac{\partial^2 T}{\partial y^2} \tag{4}$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_M \frac{\partial^2 C}{\partial y^2} + D_T \frac{\partial^2 T}{\partial y^2} \tag{5}$$

In the above system of equations,  $(u, v, w)$  are the velocity components along the  $(x, y, z)$  directions respectively,  $T$  is the fluid temperature, and  $C$  is the mass concentration.  $\mu$  represents the coefficient of fluid viscosity,  $\rho$  is the fluid density,  $\sigma$  is the constant electrical conductivity of the fluid,  $U_\infty$  is the free stream velocity, and the ratio  $\nu = (\mu/\rho)$  is the kinematic fluid viscosity. The constant parameters in the system are:  $k$  the permeability of porous material,  $C_p$  the specific heat at constant pressure,  $\kappa$  the thermal conductivity of the fluid,  $D_M$  is the molecular diffusivity and  $D_T$  is the thermal diffusivity.  $\vec{B}(x)$  is the magnetic field in the y-direction and is given by  $\vec{B}(x) = B_0/(x)^{1/2}$ .

The equation of conservation of electric charge  $\nabla \cdot J = 0$  gives  $j_y = \text{constant}$ , where  $J = (j_x, j_y, j_z)$ . This constant is assumed to be zero, since  $j_y = 0$  everywhere in the flow. The current density components  $j_x$  and  $j_z$  are obtained from the generalized Ohm's law in the following form,

$$\vec{j} = \frac{\sigma}{1 + m^2} \left( \vec{E} + \vec{V} \times \vec{B} - \frac{1}{en_e} \vec{j} \times \vec{B} \right)$$

in which  $\vec{V}$  represents the velocity vector,  $\vec{E}$  is the intensity vector of the electric field,  $\vec{B}$  is the magnetic induction vector,  $\vec{j}$  is the electric current density vector,  $m = \frac{\sigma B_0}{en_e}$  is the Hall parameter,  $\sigma$  is the electrical conductivity,  $e$  is the charge of the electron and  $n_e$  is the number density of the electron. The current density components are given by,

$$j_x = \frac{\sigma}{\alpha_e^2 + \beta_e^2} [\alpha_e (E_x - wB_0) + \beta_e (E_z + uB_0)]$$

$$j_z = \frac{\sigma}{\alpha_e^2 + \beta_e^2} [\alpha_e (E_z + uB_0) - \beta_e (E_x - wB_0)]$$

where  $\beta_e$  is the Hall current parameter,  $\beta_i$  is the Ion-slip parameter and  $\alpha_e = 1 + \beta_i \beta_e$ . In the absence of electric field ( $E_x = E_z = 0$ ), we get

$$j_x = \frac{\sigma B_0}{\alpha_e^2 + \beta_e^2} [\beta_e u - \alpha_e w]$$

$$j_z = \frac{\sigma B_0}{\alpha_e^2 + \beta_e^2} [\alpha_e u + \beta_e w]$$

The governing equations (1)-(5) become,

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

(6)

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} - \frac{\nu}{k} u - \frac{\sigma B_0^2}{\rho(\alpha_e^2 + \beta_e^2)} (\alpha_e u + \beta_e w)$$

(7)

$$u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} = \nu \frac{\partial^2 w}{\partial y^2} - \frac{\nu}{k} w + \frac{\sigma B_0^2}{\rho(\alpha_e^2 + \beta_e^2)} (\beta_e u - \alpha_e w)$$

(8)

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{\kappa}{\rho c_p} \frac{\partial^2 T}{\partial y^2}$$

(9)

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_M \frac{\partial^2 C}{\partial y^2} + D_T \frac{\partial^2 T}{\partial y^2}$$

(10)

The appropriate boundary conditions for the velocity, temperature and mass concentration are given by,

$$u = bx, v = -V_w, w = 0 \text{ at } y=0; u = w = 0 \text{ as } y \rightarrow \infty,$$

(11)

$$T = T_w \text{ at } y = 0; T \rightarrow T_\infty \text{ as } y \rightarrow \infty,$$

(12)

$$C = C_w \text{ at } y = 0; C \rightarrow C_\infty \text{ as } y \rightarrow \infty,$$

(13)

where  $b > 0$  being stretching rate of the sheet,  $V_w$  is suction/injection velocity.

### 3. Solution of the Problem

We have applied similarity technique to solve the system of equations (6)-(10) along with the boundary conditions (11)-(13). The similarity transformations are,

$$u = bxf'(\eta), v = -\sqrt{bv}f(\eta), w = \sqrt{bv}g(\eta),$$

$$\eta = \sqrt{\frac{b}{\nu}} y, \theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty}, \phi(\eta) = \frac{C - C_\infty}{(C_w - C_\infty)}$$

(14)

Introducing the above transformations in equations (6)-(10), we obtain the following system of ordinary differential equations,

$$f''' + ff'' - f'^2 - \frac{M}{(\alpha_e^2 + \beta_e^2)} (\alpha_e f' + \frac{\beta_e}{\sqrt{Re}} g) - \frac{1}{k^*} f' = 0,$$

(15)

$$g'' + fg' + \frac{M}{(\alpha_e^2 + \beta_e^2)} (\beta_e \sqrt{Re} f' - \alpha_e g) - \frac{1}{k^*} g = 0,$$

(16)

$$\theta'' + Prf\theta' = 0,$$

(17)

$$\phi'' + Scf\phi' + ScSr\theta'' = 0.$$

(18)

where  $M = \frac{\sigma B_0^2}{\rho b}$  is the Magnetic Parameter,  $Pr = \frac{\mu c_p}{\kappa}$  is the Prandtl Number,  $Sc = \frac{\nu}{D_m}$  is the Schmidt Number,  $Sr = \frac{(T_w - T_\infty) D_T}{(C_w - C_\infty) \nu}$  is the Soret number,  $k^* = kb/\nu$  is the Permeability of Porous Medium,  $Re_x = \frac{bx^2}{\nu}$  is the Reynolds Number.

In view of the similarity transformations, the boundary conditions (11)-(13) transforms to the following form,

$$f(\eta) = S, \quad f'(\eta) = 1 \text{ at } \eta = 0; \quad f'(\eta) \rightarrow 0 \text{ as } \eta \rightarrow \infty \tag{19}$$

$$g(\eta) = 0 \text{ at } \eta = 0; \quad g(\eta) \rightarrow 0 \text{ as } \eta \rightarrow \infty \tag{20}$$

$$\theta(\eta) = 1 \text{ at } \eta = 0; \quad \theta(\eta) \rightarrow 0 \text{ as } \eta \rightarrow \infty \tag{21}$$

$$\phi(\eta) = 1 \text{ at } \eta = 0; \quad \phi(\eta) \rightarrow 0 \text{ as } \eta \rightarrow \infty \tag{22}$$

Where  $S = V_w/\sqrt{b\nu}$  is the mass transfer coefficient such that  $S > 0$  indicates suction,  $S < 0$  indicates blowing at the surface.

The nonlinear coupled ordinary differential equations (15)-(18) subject to boundary conditions (19)-(22) are solved numerically using fourth order Runge-Kutta Method.

$$f' = w, \quad w' = v, \quad \theta = y, \quad \theta' = z, \quad \phi = b, \quad \phi' = e, \quad g = i, \quad g' = a.$$

$$v' = -fv + w^2 + \frac{M}{(\alpha_e^2 + \beta_e^2)} \left( \alpha_e w + \frac{\beta_e}{\sqrt{Re}} i \right) + \frac{1}{k^*} w, \tag{23}$$

$$a' = -fa - \frac{M}{(\alpha_e^2 + \beta_e^2)} (\beta_e \sqrt{Re} w - \alpha_e i) + \frac{1}{k^*} i \tag{24}$$

$$z' = -Prfz, \tag{25}$$

$$e' = -Scfe - ScSr z' \tag{26}$$

and boundary conditions becomes,

$$f(0) = S, \quad w(0) = 1, \quad i(0) = 0, \quad y(0) = 1, \quad b(0) = 1, \tag{27}$$

The above system of ODE's is reduced to a system of first order differential equations which are solved by shooting procedure using fourth order Runge Kutta method.

### 4. Parameters of Engineering Interest

The major physical quantities of engineering interest are the skin-friction coefficient  $C_f$ , the

local Nusselt number  $Nu_x$ , and the local Sherwood number  $Sh_x$  are defined respectively as follows,

Skin-friction Co-efficient:

$$C_{fx} = \frac{\tau_x}{\mu b x \sqrt{\frac{b}{\nu}}}, \text{ Here Shear stress along x-direction}$$

$$C_{fz} = \frac{\tau_z}{\mu \sqrt{b\nu} \sqrt{\frac{b}{\nu}}}, \text{ Here Shear stress along z-direction}$$

$$\tag{28}$$

$$\text{Local Nusselt Number: } Nu_x = \frac{q_w}{k \sqrt{\frac{b}{\nu}} (T_w - T_\infty)} \tag{29}$$

$$\text{Local Sherwood Number: } Sh_x = \frac{m_w}{D \sqrt{\frac{b}{\nu}} (C_w - C_\infty)} \tag{30}$$

where  $\tau_x$  and  $\tau_z$  is the skin-friction along x and z-direction respectively,  $q_w$  is the heat flux, and  $m_w$  is the mass flux at the surface.

$$\tau_x = \mu \left( \frac{\partial u}{\partial y} \right)_{y=0} = \mu b x \sqrt{\frac{b}{\nu}} f''(0)$$

$$\tau_z = \mu \left( \frac{\partial w}{\partial y} \right)_{y=0} = \mu \sqrt{b\nu} \sqrt{\frac{b}{\nu}} g'(0) \tag{31}$$

$$q_w = -k \left( \frac{\partial T}{\partial y} \right)_{y=0} = -k \sqrt{\frac{b}{\nu}} (T_w - T_\infty) \theta'(0) \tag{32}$$

$$m_w = -D \left( \frac{\partial C}{\partial y} \right)_{y=0} = -D \sqrt{\frac{b}{\nu}} (C_w - C_\infty) \phi'(0) \tag{33}$$

Substituting equations (31),(32) and (33) into the equations (28),(29) and (30), we get

$$f''(0) = C_{fx} \tag{27}$$

$$g'(0) = C_{fz}$$

$$-\theta'(0) = Nu_x$$

$$-\phi'(0) = Sh_x$$

### 5. Results and Discussion

The numerical computations are performed for several values of dimensionless parameters

involved in the equations i.e., Magnetic parameter ( $M$ ), Hall current ( $\beta_e$ ), Ion-slip parameter ( $\beta_i$ ), Schmidt Number ( $Sc$ ), Soret Number ( $Sr$ ), Suction/Blowing parameter ( $S$ ). In order to examine the computed results and understanding the behaviour of various physical parameters of the flow, numerical values for the distributions of velocity, temperature, concentration, skin friction, Nusselt number and Sherwood number are calculated by fixing various values of the non dimensional parameters of physical interest.

Figs. 2-4 show the effect of Hall current parameter  $\beta_e$  on the velocity components  $f'(\eta)$ ,  $f(\eta)$  and  $g(\eta)$ . It can be observed from Figs. 2-3 that axial velocity  $f'(\eta)$  and transverse velocity  $f(\eta)$  increase as Hall current parameter  $\beta_e$  increases. Fig. 4 depicts that cross flow velocity  $g(\eta)$  increases with increasing values of  $\beta_e$  when  $\beta_e \leq 1$ , but decreases with increasing values of Hall current parameter  $\beta_e$  greater than unity. Fig. 5 shows the effect of Hall current parameter  $\beta_e$  on concentration profiles  $\phi(\eta)$ . From this, we observe the concentration decreases with increasing values of  $\beta_e$ .

Figs. 6-8 show the influence of Hartman number on the velocity components  $f'(\eta)$ ,  $f(\eta)$  and  $g(\eta)$ . From Fig. 6-7 we observed the axial velocity and transverse velocity decrease with the increase of the magnetic field parameter values whereas the cross flow velocity  $g(\eta)$  increases with increasing values of  $M$ .

Figs. 9-11 show the effect of Ion-Slip parameter  $\beta_i$  on the velocity components  $f'(\eta)$ ,  $f(\eta)$  and  $g(\eta)$ . It is found that axial velocity  $f'(\eta)$  and transverse velocity of  $f(\eta)$  increases as  $\beta_i$  increases and cross flow velocity  $g(\eta)$  decreases with increasing values of  $\beta_i$ .

From Figs (12)-(13) we observe that the temperature  $\theta(\eta)$  and concentration profiles  $\phi(\eta)$  decrease with increasing ion-slip parameter  $\beta_i$ .

Figs. 14-16 show the influence of Permeability parameter  $k^*$  on the velocity components  $f'(\eta)$ ,

$f(\eta)$  and  $g(\eta)$ . As shown in these figures, all the velocity components are increasing with increasing values of the Permeability parameter  $k^*$ . It is expected that, an increase in the permeability of porous medium lead to a rise in the flow of the fluid through it since when the holes of the porous medium become large, the resistance of the medium may be neglected.

Fig. 17 and 18 represent the influence of Schmidt number  $Sc$  and Soret number  $Sr$  on concentration profiles  $\phi(\eta)$  respectively. It can be observed from these figures that the concentration decreases with increasing Schmidt number and increases with increasing the Soret number.

Figs. 19-23 represent velocity components  $f'(\eta)$ ,  $f(\eta)$  and  $g(\eta)$ , temperature  $\theta(\eta)$  and concentration  $\phi(\eta)$  for various values of Suction/Blowing parameter  $S$ . Here,  $S > 0$  shows the suction and  $S < 0$  shows the blowing. The axial and cross-flow velocities decrease with increasing  $S$ , whereas the transverse velocity  $f(\eta)$  increases with increasing  $S$ . The temperature and species concentration decreases as the parameter  $S$  increases.

Fig. 24 represents skin-friction coefficient  $f''(0)$  against Hall Current parameter  $\beta_e$  for various values of Ion-slip parameter  $\beta_i$ . Viscous stress acting on the surface of the plate is estimated using the skin-friction coefficient. It is clear that skin-friction increases with increasing Ion-slip  $\beta_i$  and Hall Current  $\beta_e$ .

Fig. 25 shows the Shear stress  $g'(0)$  in z-direction against Hall Current parameter  $\beta_e$  for various values of Ion-slip parameter  $\beta_i$ . Shear stress in z-direction  $g'(0)$  increases with increasing values of  $\beta_e$  when  $\beta_e \leq 1$ , but decreases with increasing values of Hall current parameter  $\beta_e$  greater than unity. Shear stress in z-direction  $g'(0)$  increases with increasing value of Ion-slip  $\beta_i$ .

Fig. 26 and 27 show that rate of heat and mass transfer decrease with the increase in the ion slip  $\beta_i$ . Increase in Ion slip increases the rate of heat transfer and decreases the rate of mass transfer.

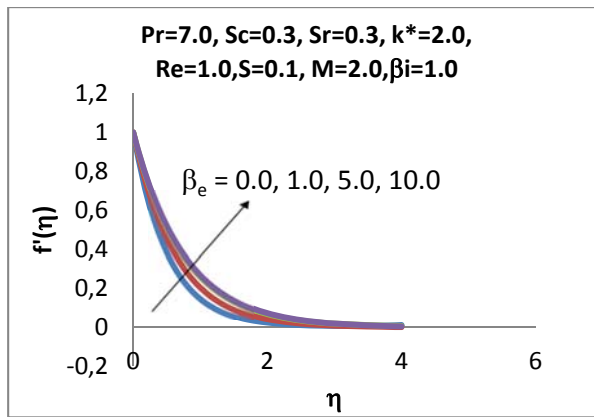


Fig. 2 Axial velocity of  $f'(\eta)$  for various values of Hall current parameter  $\beta_e$

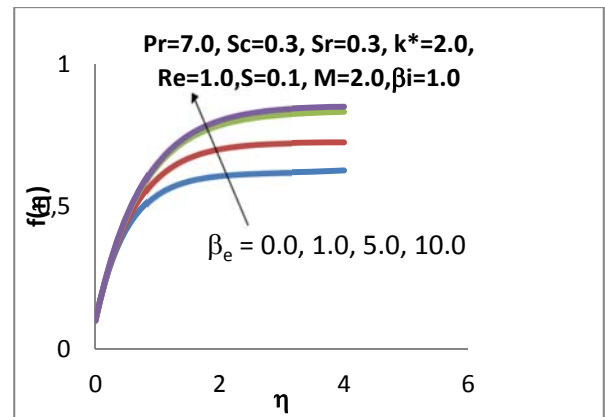


Fig. 3 Transverse velocity of  $f(\eta)$  for various values of Hall current  $\beta_e$

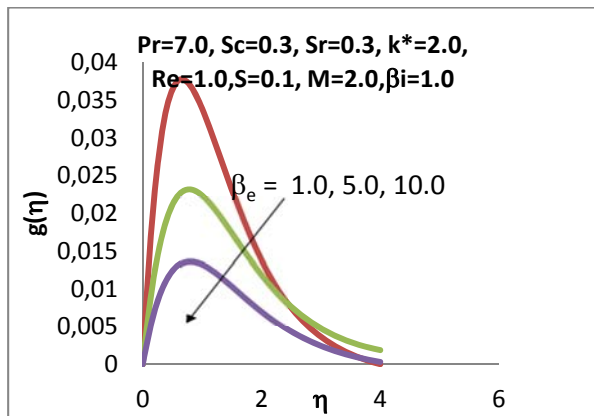


Fig. 4 Cross flow velocity of  $g(\eta)$  for various values of Hall current parameter  $\beta_e$

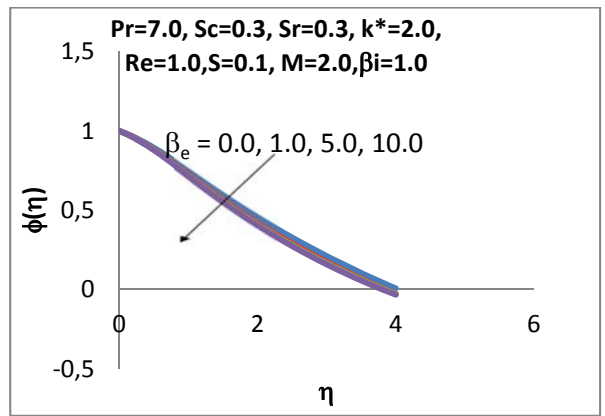


Fig. 5 Concentration profiles  $\phi(\eta)$  for various values of Hall current  $\beta_e$

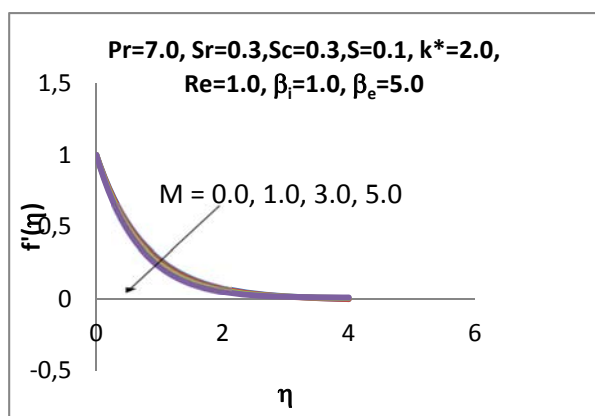


Fig. 6 Axial velocity of  $f'(\eta)$  for various values of Magnetic parameter  $M$

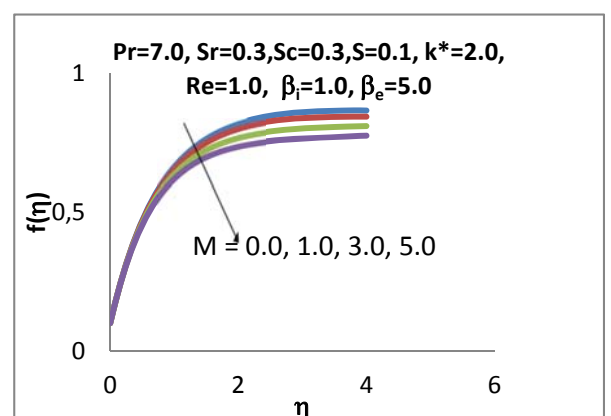


Fig. 7 Transverse velocity of  $f(\eta)$  for various values of Magnetic parameter  $M$

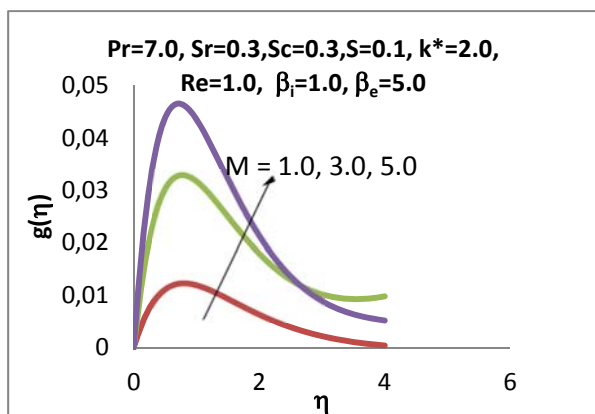


Fig. 8 Cross flow velocity of  $g(\eta)$  for various values of Magnetic parameter  $M$

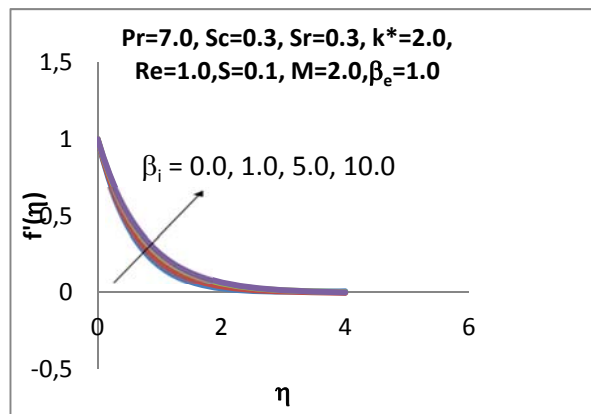


Fig. 9 Axial velocity of  $f'(\eta)$  for various values of Ion-Slip parameter  $\beta_i$

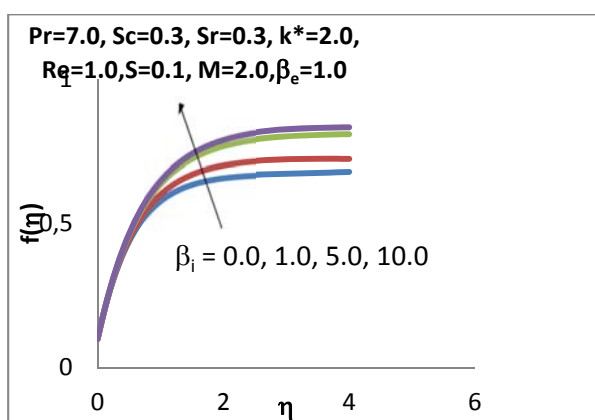


Fig. 10 Transverse velocity of  $f(\eta)$  for various values of Ion-Slip parameter  $\beta_i$

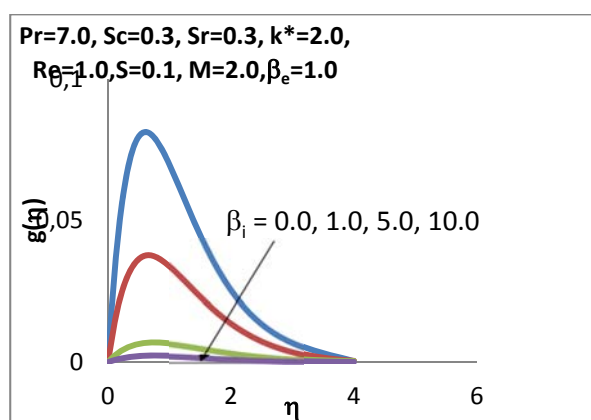


Fig. 11 Cross flow velocity of  $g(\eta)$  for various values of Ion-Slip parameter  $\beta_i$

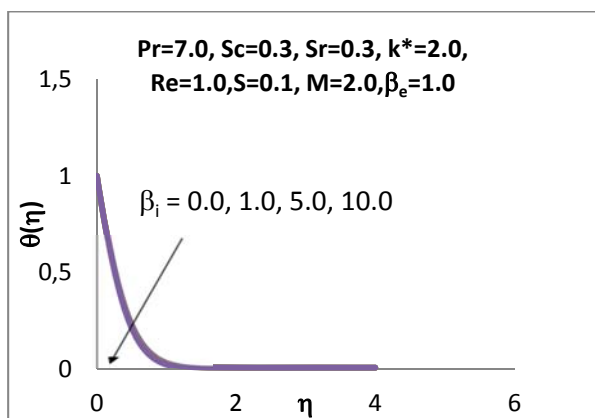


Fig. 12 Temperature profiles of  $\theta(\eta)$  for various values of Ion-Slip parameter  $\beta_i$

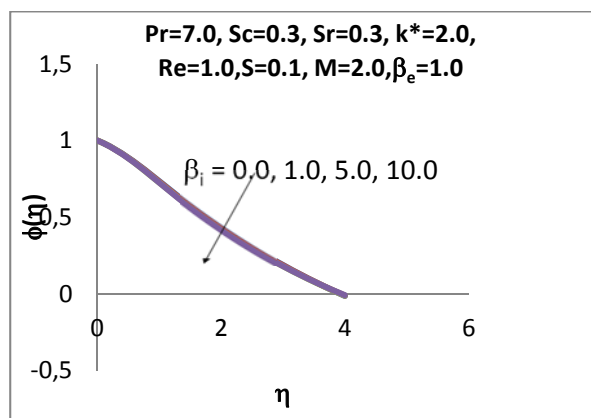


Fig. 13 Concentration profiles  $\phi(\eta)$  for various values of Ion-Slip parameter  $\beta_i$

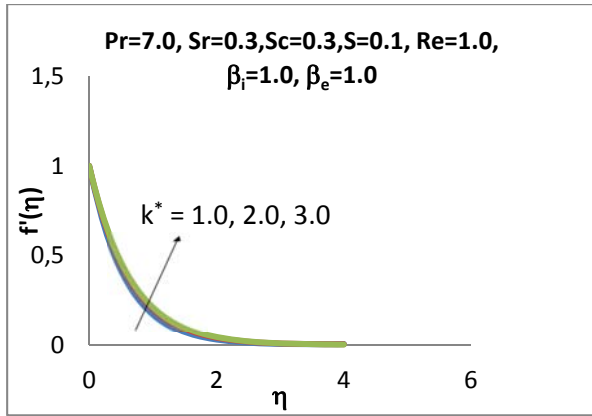


Fig. 14 Axial velocity of  $f'(\eta)$  for various values of Permeability parameter  $k^*$

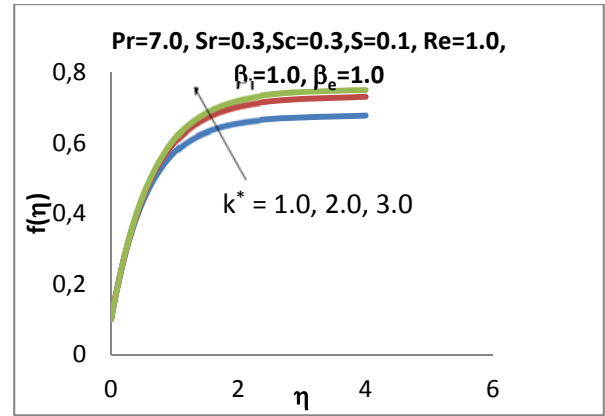


Fig. 15 Transverse velocity of  $f(\eta)$  for various values of Permeability parameter  $k^*$

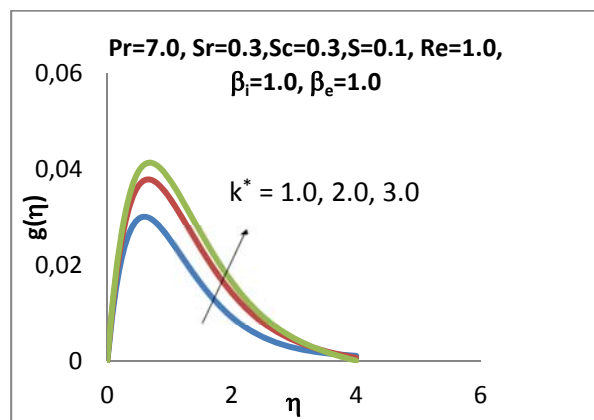


Fig. 16 Cross flow velocity of  $g(\eta)$  for various values of Permeability  $k^*$

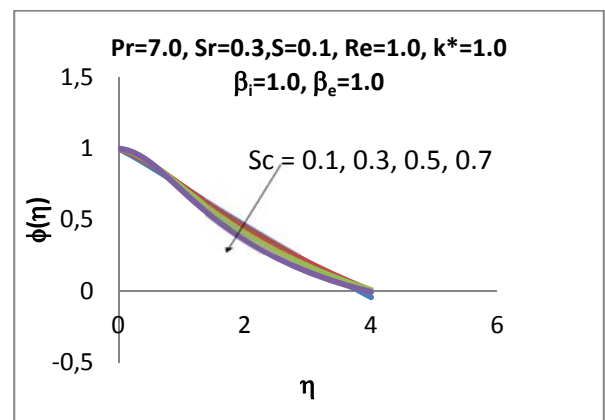


Fig. 17 Concentration profiles  $\phi(\eta)$  for various values of Schmidt number  $Sc$

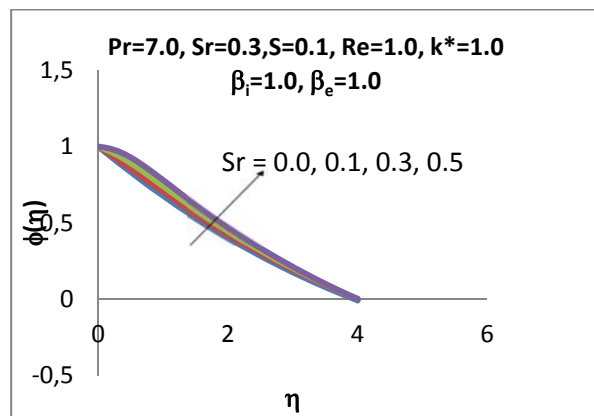


Fig. 18 Concentration profiles  $\phi(\eta)$  for various values of Soret number  $Sr$

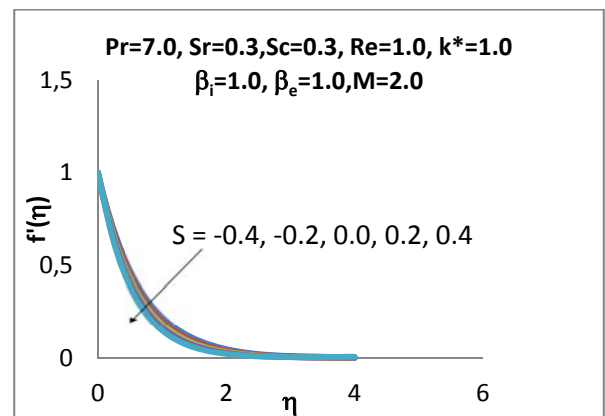


Fig. 19 Axial velocity of  $f'(\eta)$  for various values of Suction/Blowing parameter  $S$ .



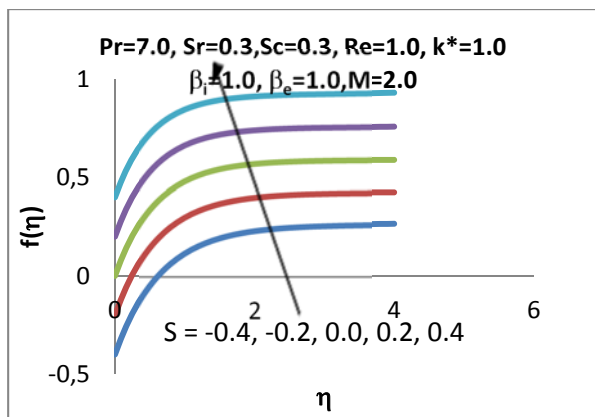


Fig. 20 Transverse velocity of  $f(\eta)$  for various values of Suction/Blowing parameter  $S$ .

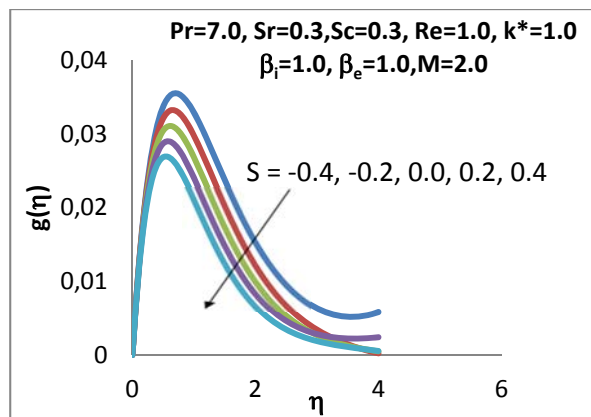


Fig. 21 Cross flow velocity of  $g(\eta)$  for various values of Suction/Blowing parameter  $S$ .

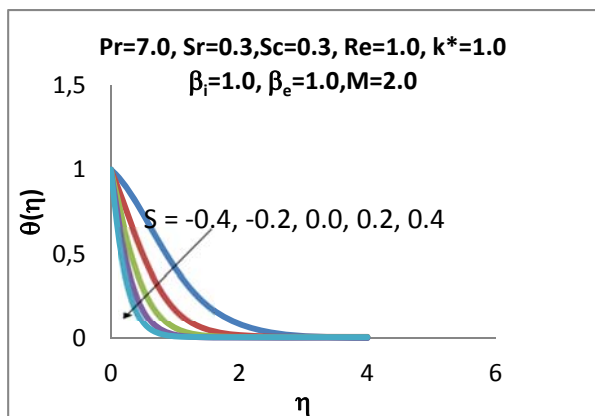


Fig. 22 Temperature profiles of  $\theta(\eta)$  for various values of Suction/Blowing parameter  $S$ .

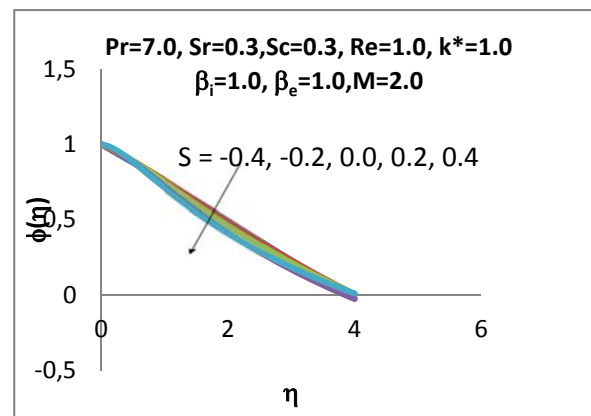


Fig. 23 Concentration profiles  $\phi(\eta)$  for various values of Suction/Blowing parameter  $S$ .

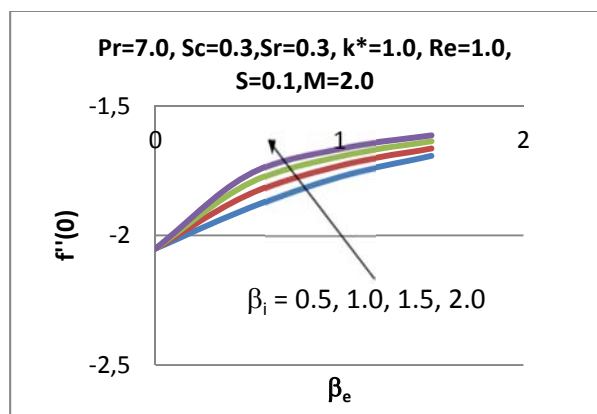


Fig. 24 Skin-friction coefficient  $f''(0)$  against Hall Current parameter  $\beta_e$  for various values of Ion-slip parameter  $\beta_i$

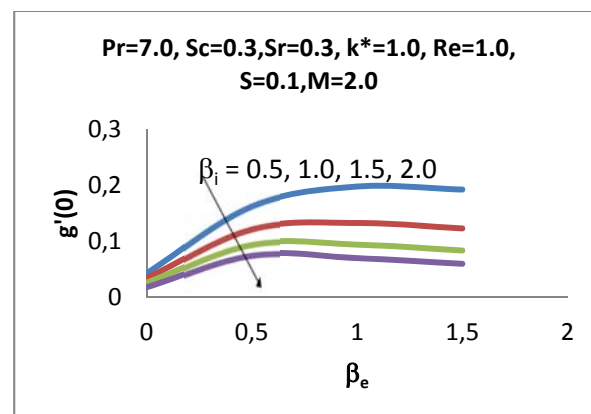


Fig. 25 Shear stress in z-direction  $g'(0)$  against Hall Current parameter  $\beta_e$  for various values of Ion-slip parameter  $\beta_i$ .

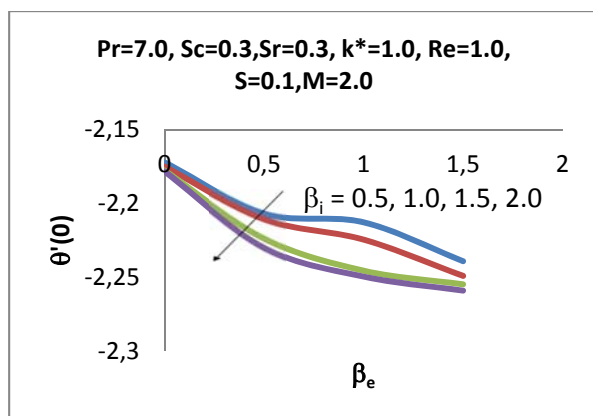


Fig. 26 Temperature gradient at the plate  $\theta'(0)$  against Hall Current parameter  $\beta_e$  for various values of Ion-slip parameter  $\beta_i$ .

## 6. Conclusion

The steady free-convective laminar flow of an incompressible viscous and electrically conducting fluid with heat and mass transfer over a porous flat plate in the presence of transverse magnetic field, Hall and ion-slip currents embedded in porous medium is investigated. The non-linear governing equations together with the boundary conditions are reduced to a system of non-linear ordinary differential equations using the similarity transformations. The system of non-linear ordinary differential equations are solved by shooting procedure using fourth order Runge-Kutta Method.

From the present study, we can conclude that all the instantaneous flow characteristics are affected by the Ion-slip parameter  $\beta_i$ . The velocity profiles  $f'(\eta), f(\eta)$  increases with increasing  $\beta_i$ . The velocity distribution i.e. cross flow velocity  $g(\eta)$  decreases with increasing values of  $\beta_i$ . Results are presented graphically to illustrate the variation shear stress, Nusselt number and Sherwood number with various parameters.

In this study the following conclusions are brought out:

- Skin-friction increases with increasing Ion-slip  $\beta_i$  and Hall Current  $\beta_e$ .
- Shear stress in z-direction  $g'(0)$  increases with increasing values of  $\beta_e$  when  $\beta_e \leq 1$ , but decreases with increasing values of Hall current parameter  $\beta_e$  greater than unity.

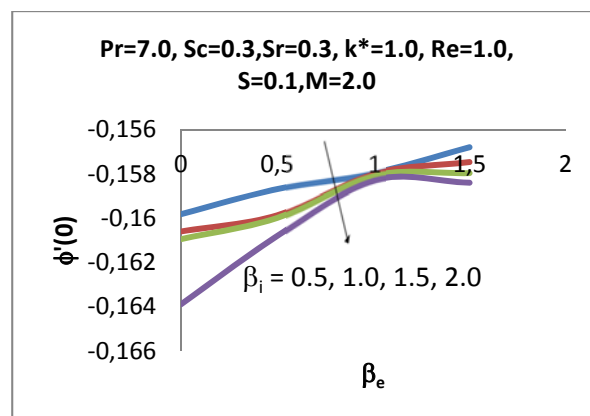


Fig. 27 Concentration gradient at the plate  $\phi'(0)$  against Hall Current parameter  $\beta_e$  for various values of Ion-slip parameter  $\beta_i$

- The rate of heat transfer  $\theta'(0)$  decreases when the increase of Ion-slip  $\beta_i$  and Hall current  $\beta_e$ .
- The rate of mass transfer  $\phi'(0)$  decreases with increase of Ion-slip  $\beta_i$  and increases with increase of Hall current  $\beta_e$ .

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